

## Temperature and substitution dependence of extremely low-energy photoemission spectra on $\text{Sm}_{1-x}\text{Eu}_x\text{B}_6$

J. Yamaguchi<sup>1</sup>, A. Sekiyama<sup>1</sup>, M. Y. Kimura<sup>1</sup>, H. Sugiyama<sup>1</sup>, S. Komori<sup>1</sup>, Y. Tomida<sup>1</sup>,  
G. Kuwahara,<sup>1</sup> T. Ito<sup>2</sup>, S. Kimura<sup>2</sup>, S. Yeo<sup>3</sup>, S.-I. Lee<sup>4</sup>, H.-D. Kim<sup>5</sup>, and S. Suga<sup>1</sup>

<sup>1</sup>Graduate school of Engineering Science, Osaka University, Toyonaka, Osaka 560-8531, Japan

<sup>2</sup>UVSOR Facility, Institute for Molecular Science, Okazaki 444-8585, Japan

<sup>3</sup>Korea Atomic Energy Research Institute, Daejeon 305-6000, Korea

<sup>4</sup>Department of Physics, Sogang University, Seoul 121-742, Korea

<sup>5</sup>Pohang Accelerator Laboratory, Pohang 790-784, Korea

$\text{SmB}_6$  and  $\text{YbB}_{12}$  have been well known as valence fluctuating (VF) Kondo semiconductors (insulators) and intensively studied because of their physical properties. The Kondo semiconductors behave as metals with localized  $f$  magnetic moments at high temperatures ( $T$ ), whereas they develop a narrow hybridization gap ( $\sim 10$  meV) at the Fermi level ( $E_F$ ) at low  $T$ . It is thought that the gap formation originates from the hybridization of the narrow  $f$  band with broad itinerant valence-bands, but the mechanism of the gap is still controversial.

In our previous hard x-ray photoemission (HAXPES) studies for the Lu substitution Kondo semiconductor alloys  $\text{Yb}_{1-x}\text{Lu}_x\text{B}_{12}$  [1], we have found from the  $T$  dependence of the Yb  $4f$  spectral analyses that the  $4f$  lattice coherence plays important roles for developing the gap. The gap for  $\text{YbB}_{12}$  is suggested to be rapidly closed by the Lu substitution of  $x = 0.125$  due to the collapse of the  $4f$  lattice coherence. On the other hand, our HAXPES study for the Eu substitution alloys  $\text{Sm}_{1-x}\text{Eu}_x\text{B}_6$  shows that the  $T$  dependence of the  $\text{Sm}^{2+}$   $4f$  spectra for  $x = 0.15$  is qualitatively similar to that for pure  $\text{YbB}_{12}$ . Thus, it is expected from our HAXPES results that a finite gap is still open for  $\text{Sm}_{0.85}\text{Eu}_{0.15}\text{B}_6$ , although the  $4f$  lattice coherence could be broken by the Eu substitution. In order to directly investigate the existence of the gap, we have performed the extremely low-energy photoemission (ELEPES) study on  $\text{Sm}_{1-x}\text{Eu}_x\text{B}_6$  ( $x = 0, 0.15, \text{ and } 0.5$ ). The ELEPES spectra for  $x = 0.15$  and  $0.5$  were measured by use of synchrotron radiation ( $h\nu = 7$  eV) at UVSOR-II BL7U and those for  $x = 0$  and  $0.5$  were measured with the Xe I ( $h\nu = 8.4$  eV) resonance line. The energy resolution was set to  $\sim 6$  meV in all measurements.

Figure 1 shows the  $T$  dependence of the ELEPES spectra near  $E_F$  for  $x = 0, 0.15, \text{ and } 0.5$ . According to the photoionization cross section [2], the observed spectra are dominated by the non- $4f$  (Sm and/or Eu  $5d$  and B  $2sp$ ) states. For  $x = 0$ , we find that a so-called leading-edge of the spectra is on the occupied side, which indicates the existence of the finite gap. In addition, the prominent peak is observed at  $\sim 15$  meV, which is comparable to that due to the magnetic excitation observed by neutron scattering

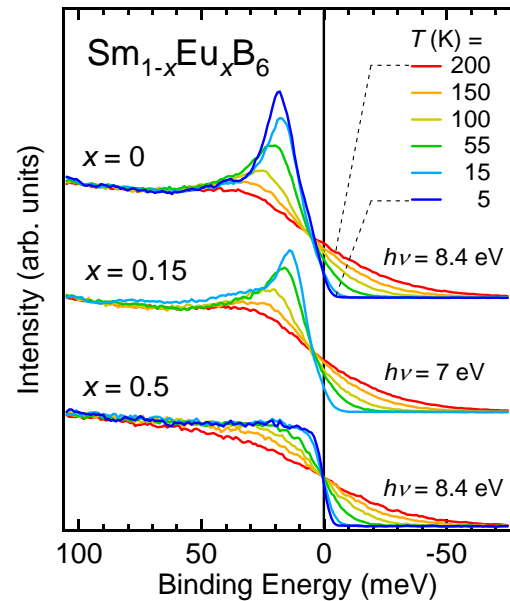


Fig. 1. Temperature dependence of the ELEPES spectra near  $E_F$  for  $\text{Sm}_{1-x}\text{Eu}_x\text{B}_6$ .

measurements for  $\text{SmB}_6$  [3]. With increasing  $T$  from 5 to 200 K, the spectral weight on  $E_F$  increases and the peak shifts toward the higher binding energy side. It should be noted that the spectra for  $x = 0.15$  show the essentially equivalent  $T$  dependence for  $x = 0$ , which indicates that  $\text{Sm}_{0.85}\text{Eu}_{0.15}\text{B}_6$  is still a Kondo semiconductor against the collapse of the  $4f$  lattice coherence. In contrast, the spectra for  $x = 0.5$  show no prominent peak and thus a typical metallic thermal behavior. These ELEPES results indicate that  $\text{SmB}_6$  is a “robust” Kondo semiconductor against a rare-earth substitution, which is significantly different from  $\text{YbB}_{12}$ .

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